

Covariant calculations of Dalitz decays of nucleon resonances and hyperons

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Plan of the talk

- **Introduction**

Formalism

- **Calculations of Dalitz decays**

- $N^* \rightarrow e^+e^-N$ decays: $N^* = \Delta(1232)_{\frac{3}{2}}^+$, $N(1520)_{\frac{3}{2}}^-$, $N(1535)_{\frac{1}{2}}^-$

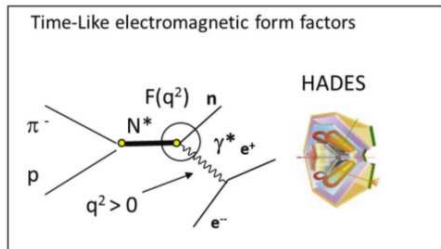
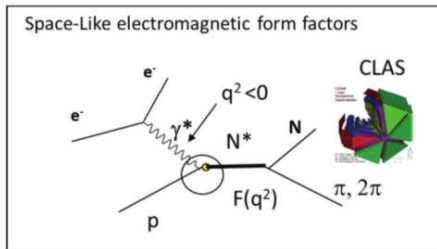
- **Hyperons:** $B' \rightarrow e^+e^-B$

Dalitz decays of decuplet baryons (B') to octet baryons (B)

- **Outlook and Conclusions**

Introduction: Spacelike and Timelike reactions

Figure: B. Ramstein, AIP Conf. Proc. 1735, 080001 (2016) [HADES]



$q^2 \leq 0$: CLAS/Jefferson Lab, MAMI,
ELSA, JLab-Hall A, MIT-BATES
 $ep \rightarrow e'N(\dots)$; $\gamma^*N \rightarrow N^*$

$q^2 > 0$: HADES,
....., PANDA
 $\pi^-p \rightarrow e^+e^-n$; $N^* \rightarrow \gamma^*N \rightarrow e^+e^-N$

Introduction: Covariant Spectator Quark Model (CSQM) †

- **Covariant Spectator Theory (BARYONS):**

Active quark off-shell; Spectator particles on-shell

Stadler, Gross and Frank PRC 56, 2396 (1998); Gross and Agbakpe PRC 73, 015203 (2006)

- Integration into quark pair d.o.f. ← time

Reduction of system to **quark-diquark** system

diquark on-shell Gross, GR and Peña PRC 77, 015202 (2008); PRD 85, 093005 (2012)

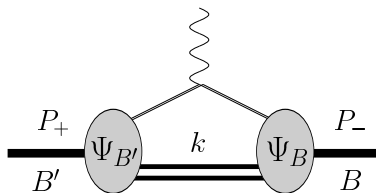
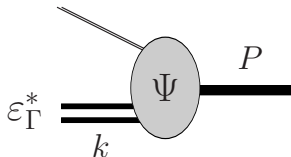
- Wave functions: relativistic generalization of $SU_S(2) \otimes SU_F(3) \otimes O(3)$ symmetries

- Radial w.f. ψ_B **determined phenomenologically** (Some physical or lattice QCD data)

- Electromagnetic interaction; relativistic impulse approximation **diquark on-shell**

$$J^\mu = 3 \sum_\Gamma \int_k \bar{\Psi}_{B'}(P_+, k) j_q^\mu \Psi_B(P_-, k)$$

- Constituent quark current $j_q^\mu(q)$ ($q = u, d, s$)
Simulate quark dressing (gluons and $q\bar{q}$ effects)



Introduction: CSQM Quark Current \oplus Transition current \dagger

● Quark current:

$$j_q^\mu(q) = j_1(q)\gamma^\mu + j_2(q)\frac{i\sigma^{\mu\nu}q_\nu}{2M_N}$$

$$j_i(q) = f_{i+}\lambda_0 + f_{i-}\lambda_3 + f_{i0}\lambda_8$$

λ_l Gell-Mann matrices; $\lambda_8 = \text{diag}(0, 0, -2)$ [strange quark]

- Functions $f_{i\ell}$ ($\ell = 0, \pm$) determined by the nucleon and decuplet lattice QCD data

F Gross, GR and MT Peña PRC 77, 015202 (2008);

GR, K Tsushima, F Gross, PRD 80, 033004 (2009)

- $f_{i\ell}(q^2)$ parametrized using VMD
Vector Meson Dominance
Include terms in m_ω , m_ρ and m_ϕ

- Quark current can be generalized to

- Lattice QCD regime
- Nuclear medium
- Timelike region (decay width Γ_v)

$$m_v \rightarrow m_v - i\Gamma_v(q)$$

GR and MT Peña, PRD 80, 013008 (2009);

GR, K Tsushima and AW Thomas, JPG 40, 015102 (2013)

Transition current \leftarrow time

$$J^\mu = J_B^\mu + J_{MC}^\mu$$

Total current =

[Bare (Valence Quark) current]

+ [Meson Cloud current]

- **Bare current:** quark model
Use calibration of quark current and radial w.f.
- **Meson cloud current:**
Educated phenomenological parametrizations

$$G_\ell(q^2) = G_\ell^B(q^2) + G_\ell^{MC}(q^2)$$

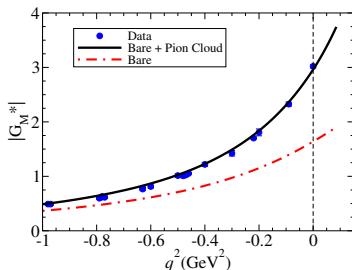
Introduction: CSQM – extension to timelike ($Q^2 = -q^2$)

Model extended from $q^2 \leq 0$ to $0 < q^2 \leq (M_{B'} - M_B)^2$

● Valence quark

Extension of quark electromagnetic form factors to $q^2 > 0$ (vector mesons)

$$\frac{m_v^2}{m_v^2 - q^2} \rightarrow \frac{m_v^2}{m_v^2 - q^2 - im_v \Gamma_v(q^2)}$$



● Meson cloud

$$\left(\frac{\Lambda_D^2}{\Lambda_D^2 - q^2} \right)^n \rightarrow \left(\frac{\Lambda_D^4}{(\Lambda_D^2 - q^2)^2 + \Lambda_D^2 [\Gamma_D(q^2)]^2} \right)^{\frac{n}{2}}$$

Effective decay width; relevant if $\Lambda_D < m_\rho$

Dalitz decay algebra ($B' \rightarrow e^+e^-B$) †

Masses: M_B initial baryon; $W =$ system γ^*B ($\approx M_{B'}$); $J^P = \frac{1}{2}^\pm, \frac{3}{2}^\pm, \dots$

Magnetic, Electric and Quadrupole Coulomb form factors; $\alpha \equiv \alpha_{em}$

G Wolf et al, NPA 517, 615 (1990); A Faessler et al, PRC 61 (2000);

M I Krivoruchenko et al, Ann. Phys. 296, 299 (2002)

$$\Gamma_{B' \rightarrow \gamma^* B}(q, W) = \alpha \mathcal{G}(M_B, W, q^2) |G_T(q^2, W)|^2 \quad [|G_T| \text{ input}]$$

$$|G_T(q^2, W)|^2 = \alpha_M |G_M(q^2, W)|^2 + \alpha_E |G_E(q^2, W)|^2 + \frac{q^2}{2W^2} |G_C(q^2, W)|^2$$

Dalitz decay rate – starting point

$$\Gamma'_{B' \rightarrow e^+e^-B}(q, W) \equiv \frac{d}{dq} \Gamma_{B' \rightarrow e^+e^-B}(q, W) = \alpha \frac{C}{q} \Gamma_{B' \rightarrow \gamma^* B}(q, W)$$

$$\Gamma_{B' \rightarrow e^+e^-B}(W) = \int_{2m_e}^{W-M_B} \Gamma'_{B' \rightarrow e^+e^-B}(q, W) dq.$$

Radiative decay: $\Gamma_{B' \rightarrow \gamma B}(W) = \Gamma_{B' \rightarrow \gamma^* B}(0, W)$; $\Gamma_{B' \rightarrow \gamma B} \equiv \Gamma_{B' \rightarrow \gamma B}(M_{B'})$

$\gamma^* N \rightarrow \Delta(1232)$ transition ($Q^2 = -q^2$)

Valence quark contribution

Dominance of magnetic dipole form factor

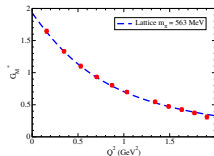
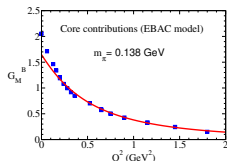
S-state wave functions $G_M^B \propto 1/Q^4$

$$G_M^B(q^2, W) = \frac{4}{3\sqrt{3}} \left[\frac{2M_N}{M_\Delta + M_N} f_{1-} + f_{2-} \right] \int_k \psi_\Delta \psi_N$$

ψ_Δ determined by EBAC data or lattice data

C Alexandrou et al Phys. Rev. D 77, 085012 (2008);

J Diaz et al, PRC 75, 015205 (2007)



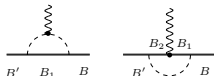
Pion (meson) cloud

Effective parametrization ($\lambda_\pi^a, \lambda_\pi^b$) based on the Cloudy Bag Model & quark counting rules (pQCD)

GR and Tsushima, PRD 88, 053002 (2013); GR, Peña Weil,

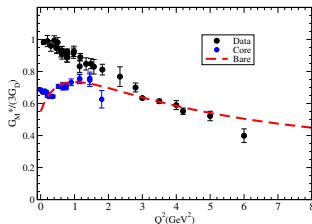
van Hees, Mosel, PRD D 93, 033004 (2016) $G_M^\pi \propto 1/Q^8$

$$G_M^\pi(q^2) = \lambda_\pi^a \left(\frac{\Lambda_\pi^2}{\Lambda_\pi^2 - q^2} \right)^2 F_\pi(q^2) + \lambda_\pi^b \left(\frac{\Lambda_D^2}{\Lambda_D^2 - q^2} \right)^4$$



(a)

(b)



Data from DESY, SLAC and CLAS (Jlab)

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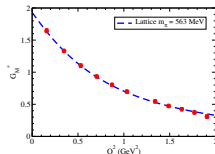
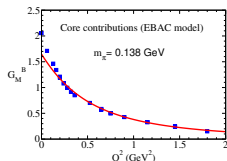
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ψ_Δ determined by EBAC data or lattice data

EBAC model: coupled-channel baryon-meson model

$\gamma N \rightarrow \pi N$ and $\gamma^* N \rightarrow \pi N \Rightarrow$ estimate core contr.

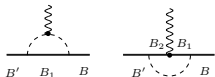


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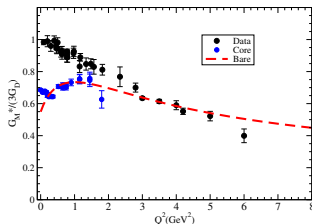
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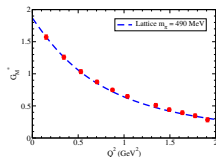
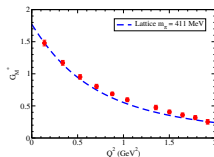
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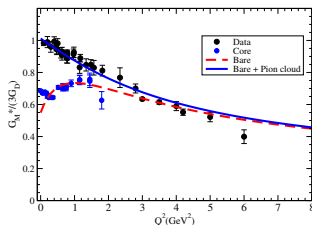
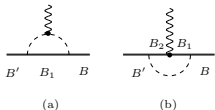
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Data from DESY, SLAC and CLAS (Jlab)

$\Delta(1232)$ Dalitz decay

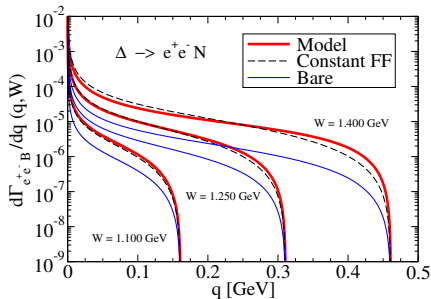
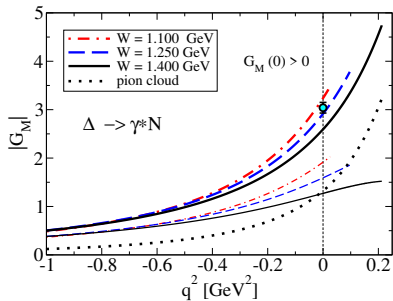
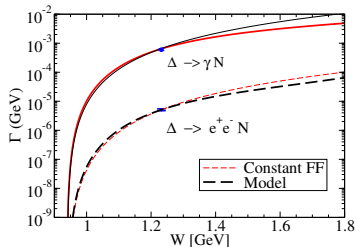
GR, MT Peña, J Weil, H van Hees and U Mosel, PRD 93, 033004 (2016)

$$\Gamma_{N^* \rightarrow \gamma^* N}(q, W) \propto$$

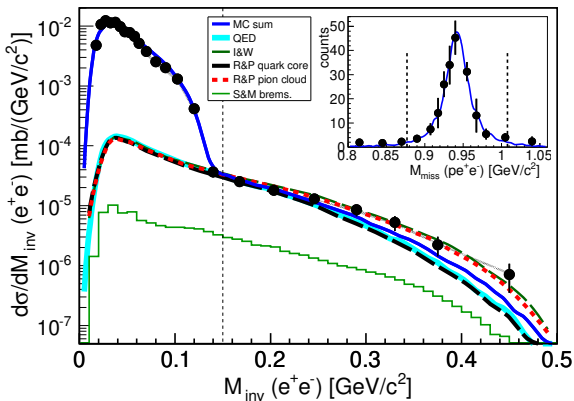
$$\left[|G_M(q^2, W)|^2 + 3|G_E(q^2, W)|^2 + \frac{q^2}{2W^2} |G_C(q^2, W)|^2 \right]$$

Simple estimate:

$$|G_T(q^2, W)|^2 \simeq |G_M(q^2, W)|^2 + \dots \quad (\dots \approx 0.2-0.6 \%)$$



$\Delta(1232)$ Dalitz decay – HADES ($pp \rightarrow p\Delta^+ \rightarrow ppe^+e^-$)



$\Delta(1232)$ Dalitz decay in proton-proton collisions at $T = 1.25$ GeV HADES, PRC 95, 065205 (2017)

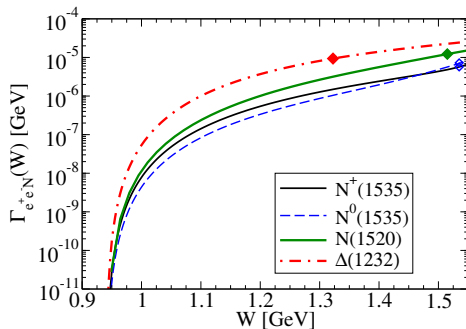
Model $\dots\dots\dots$: $\Gamma_{e^+e^-N} = 4.9$ keV

HADES: $\Gamma_{e^+e^-N} = 4.9 \pm 0.83$ keV (PDG)

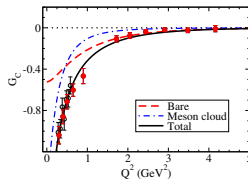
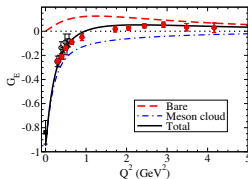
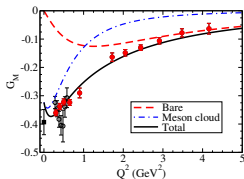
$N(1520)$ and $N(1535)$ Dalitz decay †

GR and MT Peña, PRD 95, 014003 (2017); PRD 101, 114008 (2020)

- Result dependent on isospin channel (n or p excitations)
 $N^+ \rightarrow e^+e^-p$ and $N^0 \rightarrow e^+e^-n$; TL data – information about the **neutron**
- Spacelike data dominated by **proton** data (parametrize **Meson Cloud**)



$N(1520)$ – radiative decay (MC = Meson Cloud) †

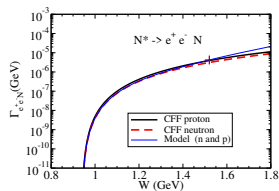
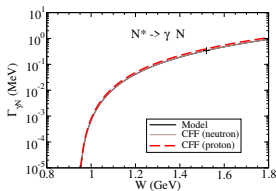


	PDG (FF)	$\Gamma_{\gamma N}$ (MeV) PDG limits	Model
p	0.43 ± 0.03	0.31–0.62	0.34
n	0.34 ± 0.03	0.30–0.64	0.34

Model: same result p and n

Dominance of MC (isovec)

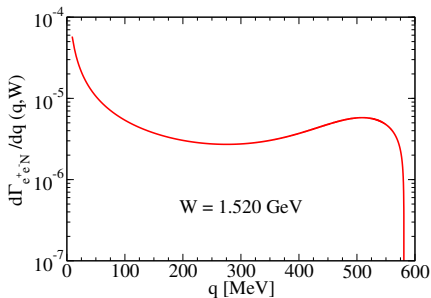
Underestimate CFF for p



SL: Valence quark contributions (Bare) dominate at large Q^2

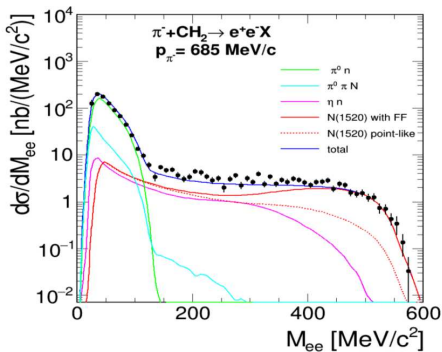
TL ($N^* \rightarrow e^+e^-N$): Model prediction larger than CFF estimates (n and p); $\Gamma_{e^+e^-B} = 4.2$ keV

$N(1520)$ Dalitz decay – HADES



Calculation: Dalitz decay rate

GR and MT Peña, PRD 95, 014003 (2017)



HADES πp $W=1.49 \text{ GeV}$,

Witold Przygoda presentation

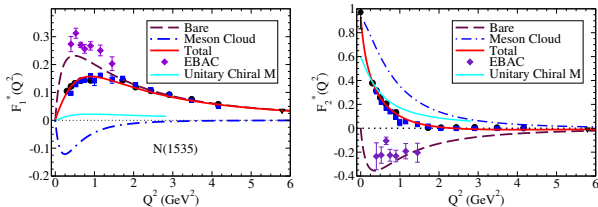
— $N(1520)$ with FF

Much larger than pointlike

Signature of q^2 -dependence of form factors $G_T(q^2, W)$

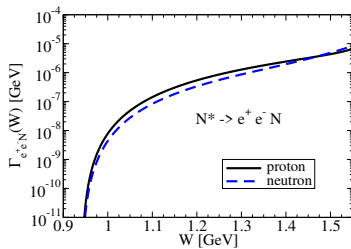
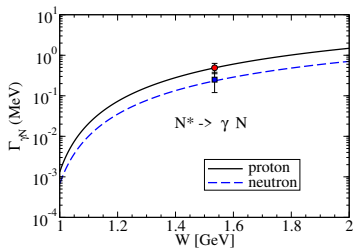
Dalitz decay measurement depend on the range of W

$N(1535)$ – radiative decay (MC = Meson Cloud) †

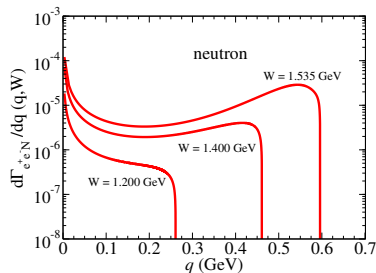
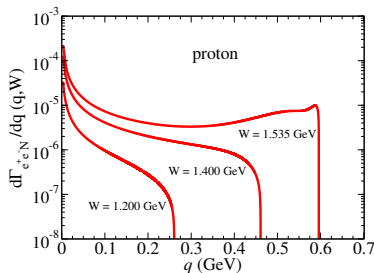


← Proton data

	$A_{1/2}(0)$ [$\text{GeV}^{-1/2}$]	$\Gamma_{\gamma N}$ [MeV]
	Data	PDG (FF) PDG limits Model
p	0.105 ± 0.015	0.49 ± 0.14 $0.19-0.53$ 0.503
n	-0.075 ± 0.020	0.25 ± 0.13 $0.013-0.44$ 0.240



$N(1535)$ Dalitz decay – HADES



↑ Calculation: Dalitz decay rate

GR and MT Peña, PRD 101, 114008 (2020)

Dependence on the isospin

q^2 -dependence related with range of W

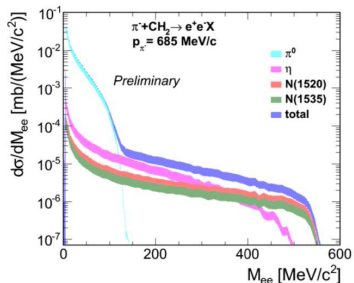
Proton: $\Gamma_{e^+e^-N} = 5.7$ keV

Neutron: $\Gamma_{e^+e^-N} = 7.2$ keV

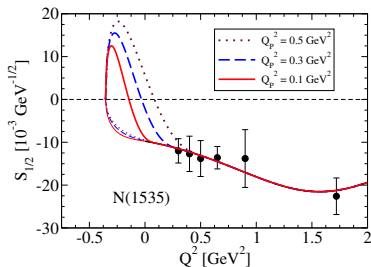
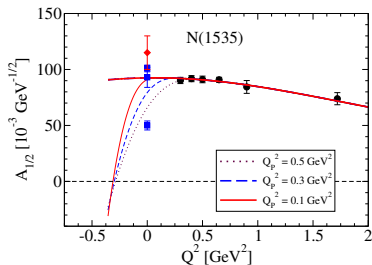
HADES, B. Ramstein, \implies

Proceedings of NSTAR2019

EPJ Web of Conf. 241, 01012 (2020)



About the $N(1535)$ data at low Q^2 (optional)



GR, PLB 759, 126 (2016); PRD 100, 114014 (2019), Devenish et al, PRD 14, 3063 (1976) — Siegart's theorem

Extension to timelike extremely sensitive to the low- Q^2 data

Ambiguities about the $\gamma^* N \rightarrow N(1535)$ data near $Q^2 = 0$ [$A_{1/2}(0) = ?$, $S_{1/2}(0) = ?$]

Estimates may change with new data

Hyperon Dalitz decays $B' \rightarrow e^+e^-B$

Extension of $\gamma^*N \rightarrow \Delta(1232)$ to **octet baryons** transitions to **decuplet baryons**

GR and K Tsushima, in preparation; GR, PRD 102, 054016 (2020); GR and K Tsushima, PRD 88, 053002 (2013)

Include also kaon cloud F_K kaon physical FF

Parameters estimated ($\Lambda_{K^a}^2, \Lambda_{K^b}^2$) by CBM expansion $G_M^{K\ell}(q^2) = G_M^{K\ell}(0) + G_M^{K\ell(1)}q^2$

$$G_M^K(q^2) = G_M^{K^a}(0) \left(\frac{\Lambda_{K^a}^2}{\Lambda_{K^a}^2 - q^2} \right)^3 F_K(q^2) + G_M^{K^b}(0) \left(\frac{\Lambda_{K^b}^4}{(\Lambda_{K^b}^2 - q^2)^2 + \Lambda_{K^b}^2 \Gamma_{K^b}^2} \right)^2$$

	$G_M^\pi(0)$	$G_M^{K^a}(0)$	$G_M^{K^b}(0)$	$G_M^K(0)$	$G_M^B(0, M_{B'})$	$\alpha_\pi(\%)$	$\alpha_K(\%)$
$\Delta \rightarrow \gamma^*N$	1.323	0.0167	0.0367	0.0534	1.633	44.0	1.8
$\Sigma^{*0} \rightarrow \gamma^*\Lambda$	1.027	0.0670	0.2768	0.3438	1.683	36.0	11.3
$\Sigma^{*+} \rightarrow \gamma^*\Sigma^+$	0.663	0.1527	0.2640	0.4167	2.094	20.9	13.1
$\Sigma^{*0} \rightarrow \gamma^*\Sigma^0$	0.270	0.1019	0.1001	0.2020	0.969	18.7	14.0
$\Sigma^{*-} \rightarrow \gamma^*\Sigma^-$	-0.124	0.0510	-0.0638	-0.0128	-0.156	42.3	4.4
$\Xi^{*0} \rightarrow \gamma^*\Xi^0$	0.308	0.1850	0.5126	0.6976	2.191	9.6	21.8
$\Xi^{*-} \rightarrow \gamma^*\Xi^-$	-0.138	0.0370	-0.1070	-0.0700	-0.168	36.7	18.6

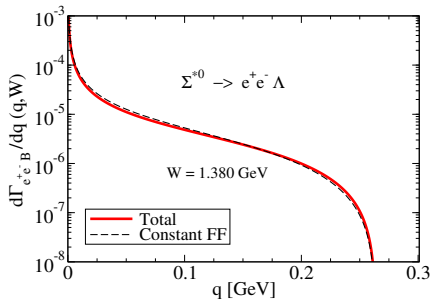
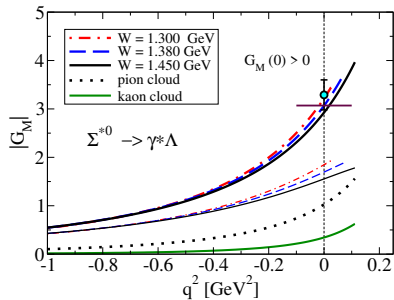
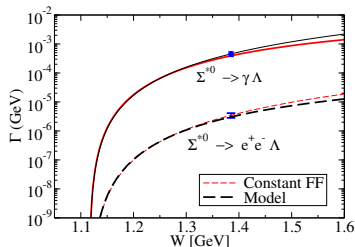
GR and K Tsushima, PRD 84, 054014 (2011); PRD 86, 114030 (2012); PRD 87, 093011 (2013);

GR, K Tsushima and AW Thomas, JPG 40, 015102 (2013)

Interesting cases: $\Sigma^{*0} \rightarrow \gamma^*\Lambda$ and $\Sigma^{*+} \rightarrow \gamma^*\Sigma^+$ ($\approx \Xi^{*0} \rightarrow \gamma^*\Xi^0$)

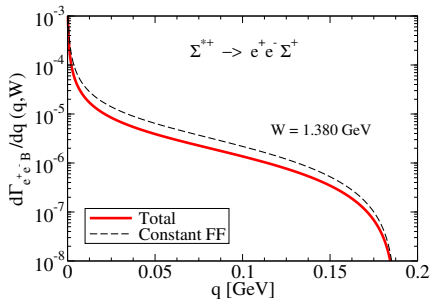
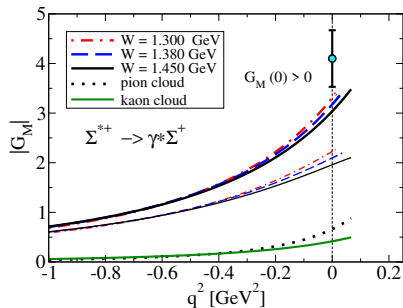
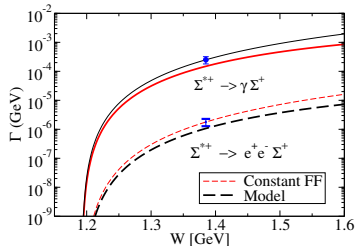
$\Sigma^{*0} \rightarrow e^+e^-\Lambda$ – Dalitz decay (preliminary)

- K increase $|G_M|$
- $|G_M(0)|_{\text{mod}} \lesssim |G_M(0)|_{\text{exp}}$
- Model \approx CFF



$\Sigma^{*+} \rightarrow e^+e^-\Sigma^+$ – Dalitz decay (preliminary)

- K increase $|G_M|$
- $|G_M(0)|_{\text{mod}} < |G_M(0)|_{\text{exp}}$
- Model $<$ CFF



Decuplet baryon decays $B' \rightarrow \gamma^* B \dagger$

Decay widths in keV $\Sigma^{*0,-}, \Xi^{*0,-}$ decays – equivalence to CFF

Decay	$\Gamma_{\gamma B}$ No K	Model	Exp	$\Gamma_{e^+e^-B}$ CFF	Model (No K)	
$\Delta \rightarrow \gamma^* N$	621	644	660±47			
$\Sigma^{*0} \rightarrow \gamma^* \Lambda$	314	392	452±86	3.47±0.65	3.04 (2.4)	5.04±0.70 ■
$\Sigma^{*+} \rightarrow \gamma^* \Sigma^+$	116	149	250±70	1.78±0.50	1.07 (0.81)	
$\Sigma^{*0} \rightarrow \gamma^* \Sigma^0$	24	31	< 1750 †	0.220	0.220 (0.16)	
$\Sigma^{*-} \rightarrow \gamma^* \Sigma^-$	1.3	1.3	< 9.5 #	0.0090	0.0091 (0.0083)	
$\Xi^{*0} \rightarrow \gamma^* \Xi^0$	111	172	†	1.253	1.265 (0.76)	
$\Xi^{*-} \rightarrow \gamma^* \Xi^-$	1.6	2.4	< 366	0.0173	0.0175 (0.012)	

■ HADES feasibility study EPJA 57, 30 (2021)

† To be measured soon? Increase strangeness production

Close to experimental limit $|G_M(0)|_{\text{mod}} \simeq 0.3$; $|G_M(0)|_{\text{exp}} < 0.8$

$$\Gamma_{\gamma B} = \frac{\alpha}{16} \frac{(M_{B'}^2 - M_B^2)^3}{M_{B'}^3 M_B^2} |G_M(0, M_{B'})|^2$$

$$\Gamma_{e^+e^-B} = \int \Gamma'_{e^+e^-B}(q, M_{B'}) dq$$

Outlook and Conclusions

Dalitz decays:

- Calculation based on **Covariant Spectator Quark Model** (Bare + Meson Cloud)
[Parametrizations sensitive to SL data ($N(1535)$) – p targets; TL data – n targets]
- Calculations with W -dependence are important to compare with data: $\Gamma_{e^+e^-B}(W)$
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 - $\Delta(1232)$ result consistent with **first Dalitz decay measurement** (HADES)
 - $\Delta(1232)$, $N(1520)$: show effects of dependence on q^2 ; $N(1535)$ under study
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 - Expected: measurements of $\Sigma^0(1385)$, $\Sigma^-(1385)$, $\Xi^0(1533)$
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 - $\Sigma^{*0} \rightarrow e^+e^-\Lambda$ and $\Sigma^{*+} \rightarrow e^+e^-\Sigma^+$: enhancement by K -cloud

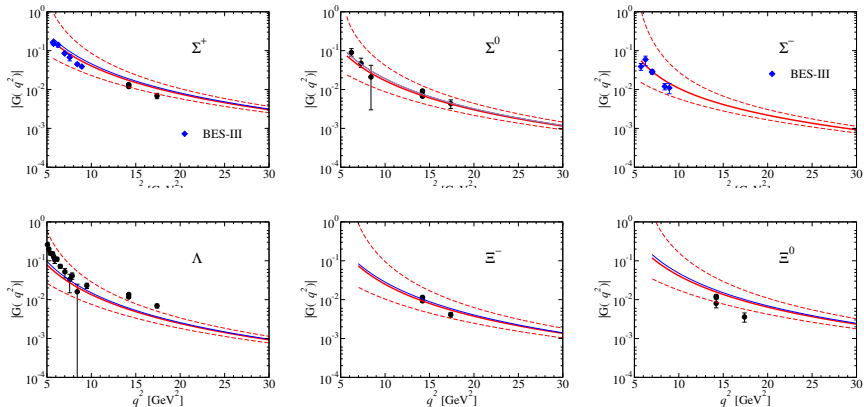
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Hyperon elastic form factors at large q^2 [Octet (Λ , Σ , Ξ), Ω^-]

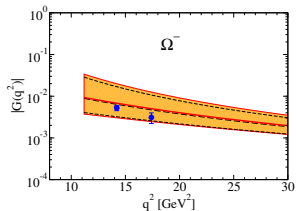
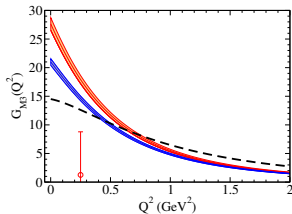
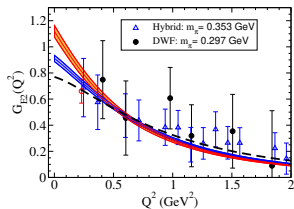
Hyperon elastic form factors at large q^2 ($e^+e^-, p\bar{p} \rightarrow B\bar{B}$)



GR, MT Peña and K Tsushima PRD 101, 014014 (2020) — Data from CLEO, BaBar, BES-III (new data Σ^-)

Ω^- elastic form factors at large q^2

GR, PRD 103, 074018 (2021)



C Alexandrou et al, PRD 82, 034504 (2010) [lattice]; S Dobbs, et al, PRD 96, 092004 (2017) [TL data]

Outlook and Conclusions

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Thank you  gilberto.ramalho2013@gmail.com